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IFUSP/P-445



A LOWER BOUND FOR THE ENERGY OF THE GROUND STATE OF BOSONS MOVING IN ONE DIMENSION

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A LOWER BOUND FOR THE ENERGY OF THE GROUND STATE OF BOSONS MOVING IN ONE DIMENSION

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SUMMARY

We show that the gound state energy of N bosons of mass m moving in one dimension is greater than $E = -\frac{m}{16h^2} \, N^2 (N-1) \, \left| \int_{-\infty}^{+\infty} V(x) \, dx \right| \quad \text{where} \quad V(x) \quad \text{is the two-body}$ potential. We conjecture that $E = -\frac{m}{24h^2} \, N \, (N^2-1) \, \left| \int_{-\infty}^{\infty} V(x) \, dx \right|^2$ provides a lower bound.

The objective of this note is to combine three known results obtaining the following theorem: the ground state of N bosons in one-dimension, interacting via a two-body potential $V(x_i-x_j)$ such that $\int_{-\infty}^{+\infty} V \, dx = -I$ (I>0), is greater than $E = -\frac{m}{16h^2} I^2 N^2 (N-1)$.

We start by recalling the three known theorems we shall use to prove the above result.

Theorem 1 - A system of N particles moving in one-dimension interacting via a globally attractive potential has always a bound state. This theorem (1) was proved by us.

Theorem 2 - Of all potentials V(x) with the same value of $I = -\int_{-\infty}^{+\infty} V(x) dx$ the 6-function potential has the lowest energy. This theorem is due to L. Spruch (2). We repeat the argument here because the proof given by Spruch is difficult to follow due to many misprints. Alternatively a more general theorem (3) was proved by J.B. Keller.

 $\underline{\text{Proof of theorem 2}}$ - Let ψ be the exact normalized ground state solution of $\ H_{\hat{0}} + V$.

So

$$E = \langle \psi | H_0 + V | \psi \rangle = \langle \psi | H_0 | \psi \rangle + \int V(x) | \psi(x) |^2 dx$$
.

Let $|\psi(x_0)|$ be the maximum value of $|\psi(x)|$. Then

 $\mathbb{E} \geq \langle \psi | \mathbb{H}_0 | \psi \rangle - | \psi (\mathbf{x}_0) |^2 \mathbb{I} = \langle \psi | \mathbb{H}_0 | \psi \rangle + \langle \psi | \nabla_{\delta} | \psi \rangle$

where $V_{\delta}(x) = -I \delta(x-x_0)$, thus concluding the proof.

 $\underline{\text{Remark}}$ - V has been assumed to be purely attractive. This assumption can be easily removed (see ref. 2).

We need now the following theorem due to Hall and $\mathsf{Post}^{(4)}$:

Theorem 3 (Hall and Post) - A lower bound for the ground state of a system of N particles is given by the ground state energy of the following two-body hamiltonian

$$H = -(N-1) \frac{\hbar^2}{2m} \Delta_{\rho_2}^2 + \frac{N}{2}(N-1) V(\sqrt{2} \rho_2)$$
, (1)

where $\rho_2 = \frac{1}{\sqrt{2}} (\vec{r}_1 - \vec{r}_2)$.

We are now ready to prove the theorem stated in the beginning of this note.

Hall-Post theorem is valid in any dimension and therefore we can say that a lower bound for the energy of $\,N$ identical bosons moving in one dimension is given by

$$E_{L} = (N-1) E_{b}$$

where $\mathbf{E}_{\mathbf{b}}$ is the ground state solution of

$$-\frac{\hbar^2}{2\mu}\frac{d^2}{dx^2}\phi(x) + \frac{N}{2}V(x)\phi(x) = E_b\phi(x) . \qquad (2)$$

But by theorem 2 we have

$$E_b > -\frac{\pi}{16\hbar^2} N^2 I^2$$
 (3)

and therefore

$$E_{L} > -\frac{m}{16\pi^{2}} (N-1) N^{2} I^{2}$$
 (4)

thus proving the result.

We now make the following:

Conjecture: A lower bound for the ground state energy of N identical particles interacting via a globally attractive potential $\left(I = -\int_{-\infty}^{+\infty} V(x) \, dx\right)$ is given by the ground state of N particles interacting via $V_{\delta} = -I \, \delta \left(x_1 - x_1\right)$.

 $\underline{\text{Quasi Proof}} \ - \ \text{Let} \quad \psi \quad \text{be the exact symmetric wave}$ function of $\ \text{H}_0\text{+V}$. Then

$$E = \langle \psi | H_0 + V | \psi \rangle = \langle \psi | H_0 | \psi \rangle + \frac{N(N-1)}{2} \int d\rho \phi(\rho) V(\rho)$$

where $\phi(\rho) = \int |\psi|^2 (x,\rho) dx$, $\rho = (x_1-x_2)$ and x stands for all the other coordinates.

Let ρ_0 be the point where $\phi(\rho,x)$ is maximum.

$$\begin{split} \mathbf{E} & \geq & <\psi \,|\, \mathbf{H}_0 \,|\, \psi > \, - \, \, \frac{\mathbf{N} \, (\mathbf{N} - \mathbf{1})}{2} \quad \, \varphi \,(\, \rho_0 \,) \quad \mathbf{I} \, = \\ \\ & = & <\psi \,|\, \mathbf{H}_0 \,|\, \psi > \, - \, \, \frac{\mathbf{N} \, (\mathbf{N} - \mathbf{1})}{2} \,\, \int \, \, \varphi \,(\, \rho \,) \,\, \mathbf{I} \,\, \delta \,(\, \rho - \rho_0 \,) \,\, \mathrm{d}\rho \,\, = \\ \\ & = & <\psi \,|\, \mathbf{H}_0 \,|\, \psi > \, - \,\, \frac{\mathbf{N} \, (\mathbf{N} - \mathbf{1})}{2} \,\, <\psi \,|\, \mathbf{I} \,\, \delta \,(\, \rho - \rho_0 \,) \,\, |\, \psi > \,\, = \\ \\ & = & <\psi \,|\, \mathbf{H}_0 \,+ \, \mathbf{U}_{\delta \, r, \rho_0} \,|\, \psi > \end{split}$$

where

$$U_{\delta,\rho_0} = -\frac{1}{2} \sum_{ij} I \delta(x_i - x_j - \rho_0)$$
.

So we have proved that the ground-state energy of $H_0 + U_{\delta,\rho_0}$ provides a lower bound for the ground-state of $H_0 + V$. We conjecture that a minimum occur for $H_0 + U_{\delta,\rho_0} = 0 = H_0 + U_{\delta}$ which is exactly solvable (5).

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